Research Article

Kinematic, Dynamic and Tribological Analysis of Push Rod Type Valve System

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Abstract

The push rod type cam and follower mechanisms are extensively used in many mechanical systems prominent among them being the internal combustion engine. It comprises of several tribo pairs: cam and follower, valve and its valve and its seat, rocker arm, tappet and its guide. The tribological performance of the various tribo interfaces affects the overall working of the mechanical system. The paper presents the kinematic and dynamic analysis of push rod type valve system for the determination of the friction between the various tribo pairs.

Keywords: Design, tribology, engine valve, cam profile, lubricant film thickness.

1. Introduction

The cam and follower mechanisms are used in variety of machines and devices. The push rod type cam and follower mechanism is shown in figure 1.



Fig.1 Push rod type valve train

2. Design of the cam profile

The determination of the cam profile is based on polynomial cam design in which equivalent two mass system of valve train is considered as shown in figure 2.



Fig.2 Dynamic model of cam and follower system

The two masses are connected by spring representing valve train linkage stiffness and hold in contact with this spring by the valve spring.

The valve end effective mass (m_1) includes all moving parts of valve train at valve end.

$$m_1 = m_v + m_{sr} + m_r + m_s /3$$

The cam end effective mass (m_2) includes all moving parts of valve train at cam end.

$$m_2 = m_t + m_p$$

Valve train stiffness (K_0) is determined by the ratio of

certain force applied to the rocker arm valve end divided by the total deflection occurring at the rocker arm valve end due to this force.

Let, y = valve lift

 y_0 = equivalent tappet lift

Considering the equilibrium of forces acting at valve end mass, we can write,

- 2

$$-K_1.y - S - K_0(y - y_0) = m_1 \cdot \frac{d^2 y}{dt^2}$$

Solving for y_0 ,

$$y_{0} = \frac{S}{K_{0}} + \frac{K_{1} + K_{0}}{K_{0}} \cdot y + \frac{m_{1}}{K_{0}} \cdot \frac{d^{2}y}{dt^{2}}$$
$$y_{0} = \frac{S}{K_{0}} + \frac{K_{1} + K_{0}}{K_{0}} \cdot y + \frac{m_{1}}{K_{0}} \cdot 36 \cdot N_{c}^{2} \cdot \frac{d^{2}y}{d\theta^{2}}$$
$$y_{0} = h_{R} + \sigma \cdot y + \delta \cdot y''$$

Thus for push rod type valve train, equivalent tappet displacement y_0 can be determined for given valve motion y. A suitable valve lift versus cam angle (θ) may be assumed and cam profile (y_c) may be developed to fulfill that motion at the desired speed. In this analysis valve lift is evaluated by using the polynomial,

$$y = L(1 + C_2 x^2 + C_4 x^4 + C_p x^p + C_q x^q + C_r x^r + C_s x^s) (1)$$

Where, $x = \frac{\theta}{\beta}$, $C_2 = -1.430657$, $C_4 = 0.1$, ${C_p} = 0.712394$, ${C_q} = -0.650403$,

$$C_r = 0.331145$$
 and $C_s = -0.06246$

 θ = Angle turned by cam

Variation of y, $\frac{dy}{d\theta}$ and $\frac{d^2y}{d\theta^2}$ versus cam angle θ is shown in figure 3.







Fig.3 Lift, velocity and acceleration variation of valve with respect to cam rotation

Then by using equation (1) equivalent tappet lift is determined. Cam lift (y_c) for each degree rotation is then determined by dividing equivalent tappet lift by rocker ratio.

$$y_c = \frac{y_0}{i} = \frac{1}{i}(h_R + \sigma.y + \delta.y')$$

Similarly,

$$y'_{c} = \frac{1}{i}(\sigma.y' + \delta.y''')$$
$$y''_{c} = \frac{1}{i}(\sigma.y'' + \delta.y^{iv})$$

3. Kinematic analysis of the cam and follower mechanism

Radius of curvature at point of contact between cam and follower is given by,

$$R = y_c''.(\frac{180}{\pi})^2 + R_b + z$$

Consequently, the velocity of point of contact on the cam surface is

$$V_c = \omega R$$

and velocity of point of contact on the follower surface is

$$V_f = \omega. y_c''. (\frac{180}{\pi})^2$$

Hence, the mean entraining velocity (V_{e}) and the sliding velocity (V_s) are given by,

$$V_{e} = \frac{1}{2} (V_{c} + V_{f})$$

55| International Journal of Advance Industrial Engineering, Vol.3, No.2 (June 2015)

$$V_s = V_c - V_f$$

4. Dynamic analysis of the cam and follower mechanism

The forces associated with the operation of valve mechanism are inertia force, spring force, forces due to dynamic deflection and damping of the components and friction between moving parts. It is assumed that the valve train is rigid and therefore the dynamic deflection and damping characteristics of the system are neglected.

In push rod valve mechanism, rocker arm plays role of transferring motion of push rod to valve. Figure 4 shows free body diagram of forces acting at rocker arm.



Fig.4 Free body diagram of rocker arm

Let, F_v be the force acting at the contact of rocker arm and valve, is given by

$$F_v = S + y.K_1 + m_v.a_v$$

where $a_v = \omega^2 . (\frac{180}{\pi}) . \frac{1}{1000} . y''$ and F_{vg} is the friction

force between valve and its guide. This interface is modeled according to fluid film lubrication assuming that the valve and its guide remain concentric and clearance space is filled with lubricant. Friction is then determined from the simple shear of lubricant film using Petroff's equation,

$$F_{vg} = \frac{\eta_0 \pi d_v l_{vg} v_v}{c_{vg}}$$

+ve sign for the friction force corresponds to the opening valve and –ve sign to the closing valve. Let, N be the reaction at rocker arm bearing and F_p be the force at the contact between push rod and rocker arm.

$$N = F_{v} + F_{p}$$

Torque equilibrium equation for rocker arm can be written as,

$$J.\ddot{\theta}_r = F_{p.}L_1 - F_{v.}L_2 - \mu_R.r_R.N$$

where μ_R is coefficient of friction in rocker arm bearing and its typical value is taken as 0.1.

Rearranging torque equilibrium equation for F_p for valve opening event,

$$F_p = \frac{J.\ddot{\theta} + F_v.L_2 + \mu_R.r_R.F_v}{L_1 - \mu_R.r_R}$$

Similarly, for valve closing event,

$$F_p = \frac{J.\ddot{\theta} + F_v.L_2 - \mu_R.r_R.F_v}{L_1 + \mu_R.r_R}$$

Let, W be the contact load between cam and tappet,

$$W = F_p + m_t a_c$$
where
$$a_c = \omega^2 \cdot (\frac{180}{\pi})^2 \cdot \frac{1}{1000} \cdot y_c''$$

$$F_{tg} = \frac{\eta_0 \pi d_t l_t v_c}{c_{tg}}$$

+ve sign for the friction force corresponds to the opening valve and -ve sign to the closing valve.

5. Lubricant film thickness

Lubricant film thickness is calculated by using Dowson and Higginson formula:

$$\frac{h}{R} = 2.65.U^{0.7}G^{0.54}W'^{-0.13}$$

Where U is dimensionless speed parameter, G is dimensionless material parameter and W' is dimensionless load parameter. Their values are given as:

$$U = \frac{\eta_0 V_e}{E'R}$$
$$G = \alpha . E'$$
$$W' = \frac{W}{E'.Rt}$$

Minimum film thickness is calculated by using above formula. The variation of film thickness along cam profile for various speeds is shown in fig. 5.





56| International Journal of Advance Industrial Engineering, Vol.3, No.2 (June 2015)

From above graph it is clear that film thickness between cam and follower increases as speed increases.

6. Friction force and torque

The cam and follower pair experiences elastohydrodynamic and boundary lubrication in one cycle. To determine friction force at the contact 'friction transition model' is used as shown in figure 6. For SAE 40W/20 oil limiting value of boundary friction coefficient is 0.12. Where λ is ratio of film thickness to composite surface roughness.



Fig.6 Friction Transition Model

According to this model, when $\lambda > 1$ friction force is determined for elastohydrodynamic lubrication and for $\lambda < 1$ friction force is determined according to boundary lubrication. The elastohydrodynamic lubrication friction force between cam and follower is,

$$F = \int_{0}^{t} \int_{-b}^{+b} \tau . dx . dy$$

Where, b is half contact length and τ is shear stress which is given by $\eta \cdot \frac{du}{dz}$. Velocity gradient $(\frac{du}{dz})$ for sliding contact of cam follower can be taken as,

$$\frac{V_c - V_f}{h}$$

Hence,

$$F = \int_{-b}^{+b} \eta \cdot \frac{V_c - V_f}{h} \cdot t \cdot dx$$

The elastic deformation between cam and follower contact is large compared with the lubricant film thickness and therefore the contact pressures and dimensions can be approximated to a dry Hertzian contact. Therefore pressure at the contact can be taken as,

$$p = p_{\text{max}} \cdot (1 - \frac{x^2}{b^2})^{\frac{1}{2}}$$

where, $p_{\text{max}} = \frac{2.W}{\pi.b.t}$
 $b = (\frac{8.W.R}{\pi.t.E'})^{\frac{1}{2}}$

According to Barus, viscosity at elevated pressure can be written as,

$$\eta = \eta_0 . e^{\alpha p}$$

Friction force is then,

$$F = \frac{\eta_0 . V_s . e^{\alpha . p_{\text{max}}}}{h} . \int_{-b}^{+b} e^{(1 - \frac{x^2}{b^2})^{1/2}} . dx$$

During boundary lubrication coefficient of friction is calculated according to λ and then friction force as,

$$F = \mu.W$$
 where, $\mu = 0.12.(1 - \lambda)$

Friction torque is then calculated for each degree of cam rotation using relation,

$$T = F.(R + y_c)$$

Using above model average friction torque is calculated for different speeds.



Fig.6 Variation of friction torque for different speeds

From above graph it is clear that as the increase in camshaft rotational frequency decreases the friction torque due to improved lubrication conditions at higher frequencies arising increased lubricant entrainment and reduced cam load at nose. At relatively low engine speeds, the cam load is dominated by valve spring compression, resulting in high friction around cam nose. At high speeds, however, inertia plays a major role as the acceleration of the tappet has negative value around the cam nose,

57| International Journal of Advance Industrial Engineering, Vol.3, No.2 (June 2015)

and it becomes more negative as engine speed increases resulting in decrease in load around the cam nose.

Conclusions

The theoretical modeling of push rod type valve system has been carried out for the determination of frictional losses. The theoretical results indicates that increase in camshaft rotational frequency decreases the friction torque due to improved lubrication conditions at higher frequencies thus resulting in reduced losses. This reduction in friction is due to the increased lubricant film thickness at higher speeds. The load around the cam nose decreases with increase in the cam rotational speed.

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